# Astrodynamics <br> (AEROOO24) 

## 3B. The Orbit in Space

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## Motivation: Space

We need means of describing orbits in three-dimensional space.


Two-body propagator


J2 propagator

## Complexity of Coordinate Systems: STK



## The Orbit in Space



Inertial frames


Coordinate systems


Coordinate types

## Importance of Inertial Frames

An inertial reference frame is defined as a system that is neither rotating nor accelerating relative to a certain reference point.

Suitable inertial frames are required for orbit description (remember that Newton's second law is to be expressed in an inertial frame).

An inertial frame is also an appropriate coordinate system for expressing positions and motions of celestial objects.

## Reference System and Reference Frame

Distinction between reference system and a reference frame:

1. A reference system is the complete specification of how a celestial coordinate system is to be formed. For instance, it defines the origin and fundamental planes (or axes) of the coordinate system.
2. A reference frame consists of a set of identifiable points on the sky along with their coordinates, which serves as the practical realization of a reference system.

## International Celestial Reference System (ICRS)

The ICRS is the reference system of the International
Astronomical Union (IAU) for high-precision astronomy.

Its origin is located at the barycenter of the solar system.

Definition of non-rotating axes:

1. The celestial pole is the Earth's north pole (or the fundamental plane is the Earth's equatorial plane).
2. The reference direction is the vernal equinox (point at which the Sun crosses the equatorial plane moving from south to north).
3. Right-handed system.

## Vernal Equinox?



The vernal equinox is the intersection of the ecliptic and equator planes, where the sun passes from the southern to the northern hemisphere (First day of spring in the northern hemisphere).

Today, the vernal equinox points in the direction of the constellation Pisces, whereas it pointed in the direction of the constellation Ram during Christ's lifetime. Why ?

## Rotation Axis: Lunisolar Precession

Because of the gravitational tidal forces of the Moon and Sun, the Earth's spin axis precesses westward around the normal to the ecliptic at a rate of $1.4^{\circ}$ century. The Earth's axis sweeps out a cone of 23.3 degrees in 26000 years.


F: dominant force on the spherical mass.
$f_{1}, f_{2}$ : forces due to the bulging sides; $f_{1}>f_{2}$, which implies a net clockwise moment.

## Rotation Axis: Lunisolar Precession



Competition between two effects:

1. Gyroscopic stiffness of the spinning Earth (maintain orientation in inertial space).
2. Gravity gradient torque (pull the equatorial bulge into the plane of the ecliptic).

## Rotation Axis: Nutation

The obliquity of the Earth varies with a maximum amplitude of $0.00025^{\circ}$ over a period of 18.6 years.

This nutation is caused by the precession of the Moon's orbital nodes. They complete a revolution in 18.6 years.


## Yet Another Disturbance: Polar Motion

Movement of Earth's rotation axis across its surface.
Difference between the instantaneous rotational axis and the conventional international origin (CIO - a conventionally defined reference axis of the pole's average location over the year 1900).

The drift, about 20 m since 1900, is partly due to motions in the Earth's core and mantle, and partly to the redistribution of water mass as the Greenland ice sheet melts.

## Yet Another Disturbance: Polar Motion



Figure 1-35. Transformation Geometry Due to Polar Motion. Accounting for polar motion takes into account the actual location of the Celestial Ephemeris Pole (CEP) over time. It moves from an ECEF system without polar motion through the CEP, to an ECEF system with polar motion using the Conventional International Origin (CIO). This correction changes the values very little, but highly accurate studies should include it. The inset plot shows the motion for the CIO from May 1986 to May 1996.

## Complicated Motion of the Earth



## Need To Specify a Date

Because the ecliptic and equatorial planes are moving, the coordinate system must have a corresponding date:
"the pole/equator and equinox of [some date]".

For ICRS, the equator and equinox are considered at the epoch J2000.0 (January 1, 2000 at 11h58m56s UTC).

## ICRS in Summary

## Quasi-equatorial coordinates at the solar system barycenter !



An object is located in the ICRS using right ascension and declination

But how to realize ICRS practically ?

## Previous Realizations: B1950 and J2000

B1950 and J2000 were considered the best realized inertial axes until the development of ICRF.

They exploit star catalogs (FK4 and FK5, respectively) which provide mean positions and proper motions for classical fundamental stars (optical measurements):

FK4 was published in 1963 and contained 1535 stars in various equinoxes from 1950 to 1975.

FK5 was an update of FK4 in 1988 with new positions for the 1535 stars.

Coord System: \begin{tabular}{ll}
\hline J2000 <br>

Prop Specific: \& | Fixed |
| :--- |
| ICRF |
| MeanDfDate |
| MeanDFEpoch |
| TrueDfDate |
| TrueDFpoch |
| B1950 |
| TEMEOFEpoch |
| TEMEOFDate |
| AlignmentAtEpoch | <br>

\hline
\end{tabular}

| 1 |  | 8 | 23.265 | +1.039 | +29 | 5 | 25.58 | -16.33 | 05 | 547.877 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2 |  | 9 | 10.695 | +6.927 | +59 | 9 | 59.18 | -19.09 | 0 | 629.765 | +6 |
| 3 |  | 9 | 24.659 | +1.196 | -45 | 44 | 50.79 | -19.11 | 0 | 652.788 | +1 |
| 4 |  | 10 | 19.257 | +0.074 | +46 | 4 | 20.21 | +0.03 | 0 | 742.779 | . |
| 5 |  | 11 | 34.437 | +0.079 | -27 | 47 | 59.12 | +1.65 | 0 | 92.265 | +0. |
| 6 |  | 11 | 44.014 | +1.412 | -35 | 7 | 59.17 | +11.86 | 0 | 911.739 | +1. |
| 7 |  | 13 | 14.154 | +0.019 | +15 | 11 | 0.80 | -1.20 | 010 | 1039.483 | +0. |
| 9 |  | 19 | 25.674 | -0.093 | -08 | 49 | 26.14 | -3.61 | 016 | 1652.829 | -0. |
| 10 |  | 20 | 4.251 | +26.779 | -64 | 52 | 29.25 | +116.39 | 017 | 1728.799 | +27 |
| 11 |  | 25 | 45.056 | +66.919 | -77 | 15 | 15.40 | +32.37 | 023 | $23 \quad 9.318$ | +68 |
| 12 |  | 26 | 17.030 | +1.933 | -42 | 18 | 21.81 | -39.57 | 023 | 2349.051 | +1 |
| 13 |  | 30 | 2.362 | +0.074 | -03 | 57 | 26.39 | -1.23 | 027 | 2729.198 | +0. |
| 14 |  | 30 | 22.661 | -0.177 | -23 | 47 | 15.72 | +1.27 | 027 | 2752.782 | -0. |
| 15 |  | 31 | 24.988 | +1.449 | -48 | 49 | 12.67 | +1.75 | 029 | 290.619 | +1 |
| 16 |  | 32 | 59.982 | +0.044 | +62 | 55 | 54.40 | -0.33 | 030 | 308.387 | . |
| 17 |  | 36 | 59.291 | +0.219 | +53 | 53 | 49.92 | -0.91 | 034 | 3410.364 | +0. |
| 19 |  | 36 | 52.858 | +0.124 | +33 | 43 | 9.63 | -0.40 | 034 | 3412.218 | +0. |
| 19 |  | 38 | 33.350 | -1.739 | +29 | 18 | 42.30 | -25.41 | 035 | 3554.458 | -1. |
| 20 |  | 39 | 19.697 | +1.060 | +30 | 51 | 39.43 | -9.15 | 036 | 3638.890 | +1 |
| 21 |  | 40 | 30.450 | +0.636 | +56 | 32 | 14.46 | -3.19 | 037 | 3739.341 | . |
| 22 |  | 43 | 35.372 | +1.637 | -17 | 59 | 11.82 | +3.25 | 041 | 414.844 | . |

Byte-by-byte description of the file: catalog

| Bytes | Format | Units | Labels | Explanations |
| :---: | :---: | :---: | :---: | :---: |
| 1- 4 | I4 | --- | FK5 | * $1 / 1670]+$ FK5 number |
| $6-7$ | I2 | h | RAh | Right ascension, hours, Equinox $=\mathrm{J} 2000$, $\mathrm{Bpoch}=\mathrm{J} 2000$ |
| 9-10 | I2 | min | RAm | Right ascension minutes (J2000.0) |
| 12-17 | F6. 3 | 3 | RAs | *Right ascension seconds (J2000.0) |
| 19-25 | F7. 3 | s/ha | PmRA | Proper motion in RA (J2000.0) |
| 27 | A1 | --- | DE- | Sign of declination (Dec) (J2000.0) |
| 28-29 | I2 | deg | DEd | Declination degrees (J2000.0) |
| 31-32 | I2 | arcmin | DEm | Declination arcminutes (J2000.0) |
| 34-38 | F5. 2 | arcsec | DEs | *Declination arcseconds (J2000.0) |
| 40-46 | F7. 2 | arcsec/ha | pmDE | Proper motion in DE (J2000.0) |
| 48-49 | I2 | h | RA1950h | Right ascension, hours <br> Equinox=B1950, Bpoch=B1950 |
| 51-52 | I2 | min | RA1950m | Right ascension minutes (B1950.0) |
| 54-59 | F6.3 | $s$ | RA1950s | *Right ascension seconds (B1950.0) |
| 61-67 | F7. 3 | s/ha | PmRA1950 | Proper motion in RA (B1950.0) |
| 69 | A1 |  | DE1950- | Sign of declination (B1950.0) |
| 70-71 | I2 | deg | DE1950d | Declination degrees (B1950.0) |
| 73-74 | I2 | arcmin | DB1950m | Declination arcminutes (B1950.0) |
| 76-80 | F5. 2 | arcsec | DE1950s | *Declination arcseconds (B1950.0) |
| 82-88 | F7. 2 | arcsec/ha | pmDE1950 | Proper motion in DE (B1950.0) |
| 90-94 | F5. 2 | a | EpRA-1900 | *Mean Epoch of observed RA |
| 96-99 | F4. 1 | ms | e_RAs | *Mean error in RA |
| 101-105 | F5. 1 | $\mathrm{ms} / \mathrm{ha}$ | e_pmRA | Mean error in pmRA |
| 107-111 | F5. 2 | a | EpDE-1900 | *Mean Epoch of observed DE |
| 113-116 | F4.1 | carcsec | e_DEs | *Mean error in Declination |
| 118-122 | F5. 1 | carcsec/ha | e_pmDE | Mean error in pmDE |
| 124-128 | F5. 2 | mag | Vmag | *V magnitude |
| 129 | A1 | --- | n_Vmag | * [VvD] Magnitude flag |
| 131-137 | A7 | --- | SpType | *Spectral type(s) |
| 139-144 | F6.3 | arcsec | plx | *?Parallax |
| 147-152 | F6.1 | km/s | RV | *?Radial velocity |
| 155-159 | A5 | --- | AGK3R | AGK3R number (Catalog $<I / 72\rangle$ ) |

## Star Catalogs: Limitations and Improvement

1. The uncertainties in the star positions of the FK5 are about 30-40 milliarcseconds over most of the sky.
2. A stellar reference frame is time-dependent because stars exhibit detectable motions.
3. Uncertainties of radio source positions are now typically less than one milliarcsecond, and often a factor of ten better.
4. Radio sources are not expected to show measurable intrinsic motion.

| 1 |  | 08 | 23.265 | +1.039 | +29 | 5 | 25.58 | -16.33 | 0 | 5 | 47.877 | +1. |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 2 |  | $\bigcirc 9$ | 10.695 | +6.927 | +59 | 9 | 59.18 | -19.09 | 0 | 6 | 29.765 | +6 |
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| 5 |  | 011 | 34.437 | +0.079 | -27 | 47 | 59.12 | +1.65 | 0 | 9 | 2.265 | +0 |
| 6 |  | 011 | 44.014 | +1.412 | -35 | 7 | 59.17 | +11.86 | 0 | 9 | 11.739 |  |
| 7 |  | 013 | 14.154 | +0.019 | +15 | 11 | 0.80 | -1.20 | 010 | 10 | 39.483 | +0 |
| 9 |  | 019 | 25.674 | -0.093 | -08 | 49 | 26.14 | -3.61 | 01 | 16 | 52.829 | -0. |
| 10 |  | 020 | 4.251 | +26.779 | -64 | 52 | 29.25 | +116.39 | 01 | 17 | 28.799 | +27 |
| 11 |  | 025 | 45.056 | +66.919 | -77 | 15 | 15.40 | +32.37 | 023 | 23 | 9.318 | +69 |
| 12 |  | 026 | 17.030 | +1.933 | -42 | 19 | 21.81 | -39.57 | 023 | 23 | 49.051 |  |
| 13 |  | 030 | 2.362 | +0.074 | -03 | 57 | 26.39 | -1.23 | 02 | 27 | 29.198 | +0 |
| 14 |  | 030 | 22.661 | -0.177 | -23 | 47 | 15.72 | +1.27 | 02 | 27 | 52.782 | -0. |
| 15 |  | 031 | 24.988 | +1.449 | -48 | 49 | 12.67 | +1.75 | 029 | 29 | 0.619 | +1 |
| 16 |  | 032 | 59.982 | +0.044 | +62 | 55 | 54.40 | -0.33 | 030 | 30 | 8.387 | +0. |
| 17 |  | 036 | 59.291 | +0.219 | +53 | 53 | 49.92 | -0.91 | 031 | 34 | 10.364 |  |
| 19 |  | 036 | 52.858 | +0.124 | +33 | 43 | 9.63 | -0.40 | 03 | 34 | 12.218 | +0. |
| 19 |  | 038 | 33.350 | -1.739 | +29 | 18 | 42.30 | -25.41 | 035 | 35 | 54.458 | -1 |
| 20 |  | 039 | 19.697 | +1.060 | +30 | 51 | 39.43 | -9.15 | 036 | 36 | 38.890 | +1 |
| 21 |  | 040 | 30.450 | +0.636 | +56 | 32 | 14.46 | -3.19 | 031 | 37 | 39.341 |  |
| 22 |  | 043 | 35.372 | +1.637 | -17 | 59 | 11.82 | +3.25 | 04 | 41 | 4.844 | +1 |

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| 9-10 | I2 | min | RAm | Right ascension minutes (J2000.0) |
| 12-17 | F6.3 | 3 | RAs | Wrigit ascerision scconcio (0̌000.0) |
| 19-25 | F7. 3 | s/ha | pmRA | Proper motion in RA (J2000.0) |
| 27 | A1 | --- | DE- | sign of declination (Dec) (J2000.0) |
| 28-29 | I2 | deg | DEd | Declination degrees (J2000.0) |
| 31-32 | I2 | arcmin | DEm | Declination arcminutes (J2000.0) |
| 34-38 | F5. 2 | arcsec | DEs | *Decilination arcseconas (J2000.0) |
| 40-46 | F7. 2 | arcsec/ha | pmDE | Proper motion in DE (J2000.0) |
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| 61-67 | F7. 3 | s/ha | PmRA1950 | Proper motion in RA (B1950.0) |
| 69 | A1 | --- | DE1950- | Sign of declination (B1950.0) |
| 70-71 | I2 | deg | DE1950d | Declination degrees (B1950.0) |
| 73-74 | I2 | arcmin | DE1950m | Declination arcminutes (B1950.0) |
| $76-80$ | F5. 2 | arcsec | DE1950s | *Declination arcseconds (B1950.0) |
| 82-88 | F7. 2 | arcsec/ha | pmDE1950 | Proper motion in DE (B1950.0) |
| 90-94 | F5. 2 | a | EpRA-1900 | *Mean Epoch of observed RA |
| 96-99 | F4.1 | ms | e_RAs | *Mean error in RA |
| 101-105 | F5. 1 | $\mathrm{ms} / \mathrm{ha}$ | e-pmRA | Mean error in pmRA |
| 107-111 | F5. 2 | a | EpDE-1900 | *Mean Epoch of observed DE |
| 113-116 | F4. 1 | carcsec | e_DEs | *Mean error in Declination |
| 118-122 | F5. 1 | carcsec/ha | e_pmDE | Mean error in pmDB |
| 124-128 | F5. 2 | mag | Vmag | *V magnitude |
| 129 | A1 | --- | n_Vmag | * [VvD] Magnitude flag |
| 131-137 | A7 | --- | SpType | *Spectral type(s) |
| 139-144 | F6. 3 | arcsec | plx | *?Parallax |
| 147-152 | F6. 1 | km/s | RV | *?Radial velocity |
| 155-159 | A5 | - | AGK3R | AGK3R number (Catalog <I/72>) |

## ICRF is the Current Realization of ICRS

Since 1998, IAU adopted the International Celestial Reference Frame (ICRF) as the standard reference frame: quasi-inertial reference frame with barely no time dependency.

It represents an improvement upon the theory behind the J2000 frame, and it is the best realization of an inertial frame constructed to date.

## Very Long Baseline Interferometry



## Further Reading on the Web Site

The Astronomical Journal, 116:516-546, 1998 July
© 1998. The American Astronomical Society. All rights reserved. Printed in U.S.A.

THE INTERNATIONAL CELESTIAL REFERENCE FRAME AS REALIZED BY VERY LONG BASELINE INTERFEROMETRY<br>C. MA<br>NASA Goddard Space Flight Center, Code 926, Greenbelt, MD 20771<br>E. F. Arias<br>Observatorio Astronómico de La Plata, Paseo del Bosque s/n, 1900 La Plata, Argentina; and Observatorio Naval Buenos Aires<br>T. M. Eubanks and A. L. Fey<br>US Naval Observatory, Code EO, 3450 Massachusetts Avenue, NW, Washington, DC 20392-5420<br>A.-M. Gontier<br>Observatoire de Paris, CNRS, URA 1125, 61 Avenue de l'Observatoire, F-75014 Paris, France<br>C. S. Jacobs and O. J. Sovers ${ }^{1}$<br>Jet Propulsion Laboratory, California Institute of Technology, 4800 Oak Grove Drive, Pasadena, CA 91109-8099<br>B. A. Archinal<br>US Naval Observatory, Code EO, 3450 Massachusetts Avenue, NW, Washington, DC 20392-5420<br>AND<br>P. Charlot ${ }^{2}$<br>Observatoire de Paris, CNRS, URA 1125, 61 Avenue de l'Observatoire, F-75014 Paris, France<br>Received 1997 December 1 ; revised 1998 March 19

TABLE 3
Coordinates of the 212 Defining Sources in the ICRF

| Designation ${ }^{\text {a }}$ | Source ${ }^{\text {b }}$ | Note ${ }^{\text {c }}$ |  |  | $\alpha(\mathrm{J} 2000.0)$ | $\delta(\mathrm{J} 2000.0)$ | $\sigma_{\alpha}$$(\mathrm{s})$ | $\begin{gathered} \sigma_{\delta} \\ (\operatorname{arcsec}) \end{gathered}$ | $C_{\alpha-\delta}$ | Epoch of Observation ${ }^{\text {d }}$ |  |  | $N_{\text {exp }}{ }^{\text {e }}$ | $N_{\text {obs }}{ }^{\text {f }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  | X | S | H |  |  |  |  |  | Mean | First | Last |  |  |
| ICRF J000557.1+382015. | $0003+380$ | ... | $\ldots$ |  | 000557.175409 | 382015.14857 | 0.000041 | 0.00051 | -0.041 | 49,087.0 | 48,720.9 | 49,554.8 | 2 | 41 |
| ICRF J001031.0 + 105829. | $0007+106$ |  |  |  | 001031.005888 | 105829.50412 | 0.000032 | 0.00068 | 0.540 | 47,938.9 | 47,288.7 | 49,690.0 | 10 | 74 |
| ICRF J001033.9 + 172418 $\ldots \ldots$. | $0007+171$ |  |  |  | 001033.990619 | 172418.76135 | 0.000021 | 0.00035 | -0.402 | 48,730.8 | 47,931.6 | 49,662.8 | 19 | 57 |
| ICRF J001331.1+405137 $\ldots \ldots$ | $0010+405$ | 2 | 1 |  | 001331.130213 | 405137.14407 | 0.000026 | 0.00034 | $-0.038$ | 49,549.6 | 48,434.7 | 49,820.5 | 7 | 219 |
| ICRF J001708.4 + 813508 $\ldots \ldots$ | $0014+813$ | $\ldots$ |  |  | 001708.474953 | 813508.13633 | 0.000121 | 0.00026 | 0.012 | 49,505.2 | 47,023.7 | 49,924.8 | 78 | 1453 |
| ICRF J004204.5 + $232001 \ldots \ldots$ | $0039+230$ |  |  |  | 004204.545183 | 232001.06129 | 0.000036 | 0.00060 | 0.090 | 48,898.1 | 48,328.5 | 49,533.8 | 3 | 44 |
| ICRF J004959.4-573827..... | 0047-579 |  |  |  | 004959.473091 | -573827.33992 | 0.000047 | 0.00053 | 0.298 | 48,697.0 | 47,626.5 | 49,407.6 | 13 | 46 |
| ICRF J011205.8 + 224438 $\ldots \ldots$. | $0109+224$ |  |  | Y | 011205.824718 | 224438.78619 | 0.000027 | 0.00049 | 0.082 | 48,733.1 | 48,434.7 | 49,736.9 | 7 | 97 |
| ICRF J012642.7 + $255901 \ldots \ldots$ | $0123+257$ |  |  |  | 012642.792631 | 255901.30079 | 0.000030 | 0.00054 | 0.167 | 48,856.4 | 48,328.5 | 49,659.8 | 4 | 71 |
| ICRF J013305.7-520003..... | 0131-522 |  |  |  | 013305.762585 | -520003.94693 | 0.000049 | 0.00081 | 0.399 | 49,039.1 | 48,162.4 | 49,895.6 | 6 | 30 |
| ICRF J013658.5 + 475129 $\ldots \ldots$ | $0133+476$ | 2 | 2 |  | 013658.594810 | 475129.10006 | 0.000026 | 0.00027 | 0.021 | 48,629.0 | 45,138.8 | 49,750.8 | 190 | 2196 |
| ICRF J013738.3-243053..... | 0135-247 |  |  |  | 013738.346378 | -24 3053.88526 | 0.000055 | 0.00042 | $-0.188$ | 48,321.8 | 47,640.2 | $49,790.7$ | 3 | 29 |
| ICRF J014125.8-092843 ..... | 0138-097 | 2 | 1 |  | 014125.832025 | -09 2843.67381 | 0.000081 | 0.00088 | 0.063 | 47,138.1 | 46,875.8 | 49,498.8 | 2 | 20 |
| ICRF J015127.1+274441 $\ldots \ldots$ | $0148+274$ |  |  |  | 015127.146149 | 274441.79365 | 0.000031 | 0.00043 | $-0.064$ | 48,963.9 | 48,328.5 | 49,659.8 | 5 | 112 |
| ICRF J015218.0 + 220707 $\ldots \ldots$. | $0149+218$ | $\ldots$ |  |  | 015218.059047 | 220707.70004 | 0.000020 | 0.00029 | -0.437 | 48,294.0 | 46,977.9 | 49,848.8 | 50 | 243 |
| ICRF J015734.9 + 744243 $\ldots \ldots$ | $0153+744$ | 4 | 3 | Y | 015734.964908 | 744243.22998 | 0.000091 | 0.00031 | 0.059 | 49,495.7 | 47,019.9 | 49,820.5 | 11 | 400 |
| ICRF J020333.3 + 723253 $\ldots \ldots$ | $0159+723$ |  |  |  | 020333.385004 | 723253.66741 | 0.000072 | 0.00031 | 0.033 | 48,800.7 | 47,011.4 | 49,667.9 | 17 | 108 |
| ICRF J020504.9 + $321230 \ldots \ldots$ | $0202+319$ |  |  |  | 020504.925371 | 321230.09560 | 0.000022 | 0.00030 | $-0.441$ | 48,017.7 | 45,466.3 | 49,736.9 | 35 | 214 |
| ICRF J021748.9 + $014449 \ldots \ldots$. | $0215+015$ | 1 | 1 |  | 021748.954740 | 014449.69909 | 0.000022 | 0.00039 | -0.215 | 49,302.1 | 48,328.5 | 49,547.8 | 5 | 133 |
| ICRF J022239.6 + 430207..... | $0219+428$ | $\ldots$ |  |  | 022239.611500 | 430207.79884 | 0.000034 | 0.00043 | $-0.098$ | 49,103.6 | 48,650.8 | 49,554.8 | 7 | 64 |
| ICRF J022256.4-344128..... | 0220-349 | ... |  |  | 022256.401625 | -34 4128.73011 | 0.000050 | 0.00044 | $-0.209$ | 48,679.5 | 47,640.2 | 49,790.7 | 4 | 35 |
| ICRF J022850.0 + $672103 \ldots \ldots$ | $0224+671$ | $\cdots$ |  |  | 022850.051459 | 672103.02926 | 0.000052 | 0.00031 | $-0.080$ | 45,097.6 | 44,090.5 | 49,600.3 | 42 | 801 |
| ICRF J022934.9-784745..... | $0230-790$ | $\cdots$ | $\cdots$ |  | 022934.946647 | -784745.60129 | 0.000149 | 0.00049 | 0.028 | 48,828.1 | 47,626.5 | 49,895.6 | 11 | 52 |
| ICRF J023838.9 + 163659 $\ldots \ldots$ | $0235+164$ | 1 | 1 |  | 023838.930108 | 163659.27471 | 0.000018 | 0.00027 | 0.090 | 47,475.7 | 44,447.0 | 49,909.6 | 194 | 2595 |
| ICRF J024229.1+110100 $\ldots \ldots$ | $0239+108$ | 2 | 2 |  | 024229.170847 | 110100.72823 | 0.000018 | 0.00030 | -0.483 | 48,582.3 | 47,511.1 | 49,662.8 | 43 | 153 |
| ICRF J025134.5 + 431515 $\ldots \ldots$ | $0248+430$ | $\ldots$ | $\ldots$ |  | 025134.536779 | 431515.82858 | 0.000027 | 0.00033 | $-0.074$ | 49,109.4 | 47,931.6 | 49,690.0 | 10 | 169 |
| ICRF J025927.0 + 074739 ..... | $0256+075$ | ... | $\ldots$ |  | 025927.076633 | 074739.64323 | 0.000021 | 0.00035 | $-0.607$ | 48,247.0 | 47,011.4 | 49,445.6 | 44 | 190 |
| ICRF J030350.6-621125..... | 0302-623 | ... |  |  | 030350.631333 | -62 1125.54983 | 0.000047 | 0.00033 | 0.129 | 49,059.2 | 48,162.4 | 49,650.8 | 15 | 97 |
| ICRF J030903.6 + 102916..... | $0306+102$ | $\ldots$ |  |  | 030903.623523 | 102916.34082 | 0.000023 | 0.00042 | $-0.804$ | 48,974.1 | 47,394.1 | 49,667.9 | 18 | 76 |
| ICRF J030956.0-605839 ..... | 0308-611 | .. | ... |  | 030956.099167 | -6058 39.05628 | 0.000038 | 0.00029 | 0.037 | 49,029.5 | 47,626.5 | 49,895.6 | 79 | 738 |
| ICRF J031301.9 + 412001 $\ldots$... | $0309+411$ | $\ldots$ | $\ldots$ | Y | 031301.962129 | 412001.18353 | 0.000026 | 0.00031 | $-0.321$ | 48,371.0 | 47,165.8 | 49,848.8 | 29 | 127 |
| ICRF J034506.4 + $145349 \ldots \ldots$ | $0342+147$ | . | $\cdots$ |  | 034506.416546 | 145349.55818 | 0.000021 | 0.00032 | $-0.622$ | 48,809.6 | 47,394.1 | 49,445.6 | 23 | 177 |
| ICRF J040305.5 + $260001 \ldots \ldots$ | $0400+258$ | 3 | 2 | Y | 040305.586048 | 260001.50274 | 0.000020 | 0.00030 | -0.127 | 48,990.5 | 47,005.8 | 49,820.5 | 37 | 397 |
| ICRF J040922.0 + 121739 $\ldots \ldots$ | $0406+121$ | 2 | 1 |  | 040922.008740 | 121739.84750 | 0.000021 | 0.00033 | -0.704 | 48,399.2 | 46,977.9 | 49,565.9 | 28 | 149 |
| ICRF J041636.5-185108..... | 0414-189 | $\ldots$ | $\ldots$ |  | 041636.544466 | -185108.34012 | 0.000051 | 0.00048 | $-0.078$ | 47,814.6 | 46,840.8 | 49,790.7 | 3 | 31 |
| ICRF J042442.2-375620 $\ldots \ldots$ | 0422-380 | $\cdots$ | $\cdots$ |  | 042442.243727 | -375620.78423 | 0.000033 | 0.00119 | 0.251 | 49,081.7 | 48,162.4 | 49,750.8 | 11 | 60 |
| ICRF J042446.8+003606..... | $0422+004$ | 2 | 1 |  | 042446.842052 | 003606.32983 | 0.000020 | 0.00063 | 0.038 | 48,938.2 | 45,997.8 | 49,820.5 | 11 | 245 |
| ICRF J042636.6+051819 $\ldots \ldots$ | $0423+051$ | $\ldots$ | $\ldots$ |  | 042636.604102 | 051819.87204 | 0.000031 | 0.00087 | 0.101 | 48,977.3 | 48,194.7 | 49,667.9 | 9 | 64 |
| ICRF J042840.4-375619..... | 0426-380 | $\ldots$ | $\ldots$ |  | 042840.424306 | -375619.58031 | 0.000036 | 0.00047 | 0.011 | 48,125.7 | 47,640.2 | 49,692.6 | 5 | 39 |
| ICRF J043900.8-452222..... | 0437-454 | $\ldots$ | $\ldots$ |  | 043900.854714 | -4522 22.56260 | 0.000057 | 0.00078 | $-0.123$ | 49,443.5 | 48,766.9 | 49,895.6 | 7 | 32 |
| ICRF J044238.6-001743..... | 0440-003 | 1 | 1 |  | 044238.660762 | -0017 43.41910 | 0.000025 | 0.00064 | 0.262 | 47,735.2 | 47,011.4 | 49,576.9 | 15 | 111 |
| ICRF J044907.6 + 112128 $\ldots \ldots$ | $0446+112$ | $\ldots$ | $\ldots$ |  | 044907.671119 | 112128.59662 | 0.000024 | 0.00051 | $-0.143$ | 49,312.0 | 47,394.1 | 49,854.8 | 5 | 32 |
| ICRF J045005.4-810102..... | 0454-810 | ... | ... |  | 045005.440195 | -810102.23146 | 0.000137 | 0.00032 | $-0.005$ | 48,784.2 | 47,626.5 | 49,895.6 | 18 | 148 |
| ICRF J045952.0 + 022931..... | $0457+024$ | $\ldots$ | $\ldots$ |  | 045952.050664 | 022931.17631 | 0.000019 | 0.00032 | 0.062 | 48,993.4 | 47,005.8 | 49,750.8 | 36 | 394 |
| ICRF J050145.2 + 135607 . | $0458+138$ | 2 | 2 |  | 050145.270840 | 135607.22063 | 0.000037 | 0.00064 | $-0.770$ | 48,830.7 | 47,394.1 | 49,848.8 | 13 | 20 |
| ICRF J050523.1+045942 $\ldots \ldots$ | $0502+049$ | $\ldots$ | $\ldots$ |  | 050523.184723 | 045942.72448 | 0.000037 | 0.00060 | $-0.584$ | 48,897.7 | 47,394.1 | 49,667.9 | 6 | 28 |
| ICRF J050643.9-610940 $\ldots \ldots$. | 0506-612 | $\ldots$ |  |  | 050643.988739 | -6109 40.99328 | 0.000047 | 0.00035 | 0.145 | 48,760.5 | 48,110.9 | 49,594.7 | 16 | 69 |
| ICRF J050842.3+843204 $\ldots \ldots$. | $0454+844$ | $\ldots$ | $\ldots$ |  | 050842.363503 | 843204.54402 | 0.000194 | 0.00028 | -0.046 | 48,674.7 | 46,977.9 | 49,611.9 | 42 | 250 |
| ICRF J051002.3 + 180041 $\ldots \ldots$. | $0507+179$ | 2 | 2 |  | 051002.369122 | 180041.58171 | 0.000020 | 0.00030 | $-0.396$ | 49,401.9 | 47,605.1 | 49,820.5 | 24 | 339 |
| ICRF J051644.9-620705..... | 0516-621 | $\ldots$ | $\ldots$ |  | 051644.926178 | -620705.38930 | 0.000048 | 0.00042 | 0.202 | 49,455.4 | 48,749.6 | 49,895.6 | 9 | 56 |

## Formal Definition of ICRS

It is defined by the measured positions of 212 extragalactic sources (mainly quasars).

1. Its origin is located at the barycenter of the solar system through appropriate modeling of VLBI observations in the framework of general relativity.
2. Its pole is in the direction defined by the conventional IAU models for precession (Lieske et al. 1977) and nutation (Seidelmann 1982).
3. Its origin of right ascensions was implicitly defined by fixing the right ascension of the radio source 3C273B to FK5 J2000 value.

## 3. The Orbit in Space



Coordinate systems

## Coordinate Systems

Now that we have defined an inertial reference frame, other reference frames can be defined according to the needs of the considered application.

Coordinate transformations between two reference frames involve rotation and translation.

What are the possibilities for a satellite in Earth orbit?

## Geocentric — Inertial (ECI)

A geocentric-equatorial system is clearly convenient.

The geocentric celestial reference frame (GCRF) is the geocentric counterpart of the ICRF and is the standard inertial coordinate system for the Earth.

## Geocentric - Fixed (ECEF)

Origin at the Earth's center.
$\Rightarrow z$-axis is parallel to Earth's rotation vector.
$\Rightarrow x$-axis passes through the Greenwich meridian.
$\Rightarrow y$-axis: right-handed set.

For ground tracks and force computation.

ECEF


## ECEF-ECI Transformation

It includes precession, nutation, and rotation effects, as well as pole wander and frame corrections.


## ECEF-ECI Transformation



Figure 3-29. Classical Transformation. This figure depicts the transformation of a state vector in the body fixed (ITRF) frame to the inertial (FK5) frame. This two-way conversion is necessary for many orbit determination problems. The clear ellipses show the intermediate frames.

Vallado, Fundamental of Astrodynamics and Applications, Kluwer, 2001.

## ECEF-ECI Transformation

Simplified transformation

$$
\begin{aligned}
& \omega_{\oplus}=0.000,072,921,158,553,0 \mathrm{rad} / \mathrm{s} \\
& \theta_{\mathrm{GMST}, 2000}=1.74476716333061 \mathrm{rad} \\
& \theta_{\mathrm{GMST}}=\theta_{\mathrm{GMST}, 2000}+\omega_{\oplus} \times 86400 \times(t+0.5) \mathrm{rad} \\
& \left(\begin{array}{l}
\mathrm{x} \\
\mathrm{y} \\
\mathrm{z}
\end{array}\right)_{E C I}=\left[\begin{array}{ccc}
\cos \left(\theta_{\mathrm{GMST}}\right) & -\sin \left(\theta_{\mathrm{GMST}}\right) & 0 \\
\sin \left(\theta_{\mathrm{GMST}}\right) & \cos \left(\theta_{\mathrm{GMST}}\right) & 0 \\
0 & 0 & 1
\end{array}\right]\left(\begin{array}{l}
x \\
y \\
z
\end{array}\right)_{E C E F}
\end{aligned}
$$

Precession, nutation, polar motion ignored


| Control |
| :---: |
| (o) None <br> () Cross-Section <br> Attitude |



## Integrator

- ODE113
( ) RK8(7)
- RK8
$\square$ Download Data
ECl to ECEF
- $\downarrow$ Precession

V Nutation
V Polar Wandering
Simplified


Density Model
Harris-Priester

- Jacchia 71
(O) Jacchia-Roberts
$\square$ Measured data
Density Parameters
Harris-Priester coeff.
DailyF10.7
Averaged F10.7

| 0 |
| :---: |
| 155 |
| 155 |
| 3 |

Spacecraft Properties


## Yet More Coordinate Systems !

Satellite coordinate system

Perifocal coordinate system

Heliocentric coordinate system

Non-singular elements

## For ADCS

Natural frame for an orbit (z is zero)

For interplanetary missions

For particular orbits

## 3. The Orbit in Space



## Cartesian and Spherical

1. Cartesian: for computations
2. Spherical: azimuth and elevation (for ground station) right ascension and declination (for astronomers)


## Cartesian $\leftrightarrow$ Spherical

$$
\mathbf{r}=X \hat{\mathbf{I}}+Y \hat{\mathbf{J}}+Z \hat{\mathbf{K}}=r \hat{\mathbf{u}}_{r}
$$

$\hat{\mathbf{u}}_{r}=\cos \delta \cos \alpha \hat{\mathbf{I}}+\cos \delta \sin \alpha \hat{\mathbf{J}}+\sin \delta \hat{\mathbf{K}}$

## Orbitron



## Orbitron: Close-Up



## Orbital (Keplerian) Elements

## For interpretation

$\mathbf{r}$ and $\mathbf{v}$ do not directly yield much information about the orbit. We cannot even infer from them what type of conic the orbit represents !

Another set of six variables, which is much more descriptive of the orbit, is needed.

## 6 Orbital (Keplerian) Elements

1. e: shape of the orbit
2. a: size of the orbit
3. $i$ : orients the orbital plane with respect to the ecliptic plane
4. $\Omega$ : longitude of the intersection of the orbital and ecliptic planes
5. $\omega$ : orients the semi-major axis with respect to the ascending node
6. $v$ : orients the celestial body in space
position of the satellite on the ellipse
orientation of the ellipse
position of the satellite



| Control |
| :---: |
| (o) None <br> () Cross-Section <br> Attitude |



## Integrator

- ODE113
( ) RK8(7)
- RK8
$\square$ Download Data
ECl to ECEF
- $\downarrow$ Precession

V Nutation
V Polar Wandering
Simplified


Density Model
Harris-Priester

- Jacchia 71
(O) Jacchia-Roberts
$\square$ Measured data
Density Parameters
Harris-Priester coeff.
DailyF10.7
Averaged F10.7

| 0 |
| :---: |
| 155 |
| 155 |
| 3 |

Spacecraft Properties


## Orbital Elements a,e,i, $\Omega, \omega, \theta$ from r,v ?



## e and a from the 2-body Problem

$$
\mu \boldsymbol{e}=\boldsymbol{v} \times \boldsymbol{h}-\mu \frac{\boldsymbol{r}}{r}
$$

$$
v=\sqrt{\mu\left(\frac{2}{r}-\frac{1}{a}\right)}
$$



$$
e=\left\|\frac{\boldsymbol{v} \times(\boldsymbol{r} \times \boldsymbol{v})}{\mu}-\frac{\boldsymbol{r}}{r}\right\|
$$



$$
a=\frac{r}{2-\frac{r v^{2}}{\mu}}
$$

$$
r=\|\boldsymbol{r}\|, v=\|v\|
$$

## Inclination

Angle between the orbital and equatorial planes:


$$
i=\cos ^{-1}\left(\frac{(\boldsymbol{r} \times \boldsymbol{v}) \cdot \widehat{\boldsymbol{K}}}{\|\boldsymbol{r} \times \boldsymbol{v}\|}\right)
$$

## Longitude $\Omega$

Angle between the nodal vector $\mathbf{n}$ and the vernal equinox:

$$
\cos \Omega=\frac{\boldsymbol{n} . \hat{\boldsymbol{I}}}{\|\boldsymbol{n}\|}
$$

The nodal vector $\mathbf{n}$ is in the orbital and equatorial planes:

$$
\boldsymbol{n}=\widehat{\boldsymbol{K}} \times \frac{\boldsymbol{h}}{h}
$$



$$
\Omega=\cos ^{-1} \frac{\boldsymbol{n} \cdot \widehat{\boldsymbol{I}}}{\|\boldsymbol{n}\|}=\cos ^{-1}\left(\frac{\left(\widehat{\boldsymbol{K}} \times \frac{\boldsymbol{r} \times \boldsymbol{v}}{\|\boldsymbol{r} \times \boldsymbol{v}\|}\right) \cdot \hat{\boldsymbol{I}}}{\left\|\widehat{\boldsymbol{K}} \times \frac{\boldsymbol{r} \times \boldsymbol{v}}{\|\boldsymbol{r} \times \boldsymbol{v}\|}\right\|}\right)
$$

n. $\widehat{J} \geq \mathbf{0}$

$$
\Omega=360^{\circ}-\Omega
$$

n. $\widehat{J}<\mathbf{0}$

## Argument of Perigee

Angle between the nodal and eccentricity vectors:

$$
\cos \omega=\frac{\boldsymbol{e} . \boldsymbol{n}}{\|\boldsymbol{e}\|\|\boldsymbol{n}\|}
$$

$$
\sum n=\widehat{\boldsymbol{K}} \times \frac{\boldsymbol{h}}{h}, \boldsymbol{e}=\frac{\boldsymbol{v} \times(\boldsymbol{r} \times \boldsymbol{v})}{\mu}-\frac{\boldsymbol{r}}{r}
$$

$$
\begin{array}{rr}
\omega=\cos ^{-1}\left(\frac{\left(\widehat{\boldsymbol{K}} \times \frac{\boldsymbol{r} \times \boldsymbol{v}}{\|\boldsymbol{r} \times \boldsymbol{v}\|}\right) \cdot\left(\frac{\boldsymbol{v} \times(\boldsymbol{r} \times \boldsymbol{v})}{\mu}-\frac{\boldsymbol{r}}{r}\right)}{\left\|\widehat{\boldsymbol{K}} \times \frac{\boldsymbol{r} \times \boldsymbol{v}}{\|\boldsymbol{r} \times \boldsymbol{v}\|}\right\|\left\|\frac{\boldsymbol{v} \times(\boldsymbol{r} \times \boldsymbol{v})}{\mu}-\frac{\boldsymbol{r}}{r}\right\|}\right) & \text { e. } \widehat{\boldsymbol{K}} \geq \mathbf{0} \\
\omega=360^{\circ}-\omega & \text { e. } \widehat{\boldsymbol{K}}<\mathbf{0}
\end{array}
$$

## True Anomaly

Angle between the position and eccentricity vectors

$$
\begin{gathered}
\cos \theta=\frac{\boldsymbol{r} \cdot \boldsymbol{e}}{r\|\boldsymbol{e}\|} \\
\theta=\cos ^{-1}\left(\frac{\boldsymbol{r} \cdot\left(\frac{\boldsymbol{v} \times(\boldsymbol{r} \times \boldsymbol{v})}{\mu}-\frac{\boldsymbol{r}}{r}\right)}{r\left\|\frac{\boldsymbol{v} \times(\boldsymbol{r} \times \boldsymbol{v})}{\mu}-\frac{\boldsymbol{r}}{r}\right\|}\right) \\
\theta=\mathbf{r} \cdot \boldsymbol{v} \geq \mathbf{0} \\
\theta=360^{\circ}-\theta \\
\boldsymbol{r} \cdot \boldsymbol{v}<\mathbf{0}
\end{gathered}
$$

## r,v from a,e,i, $\Omega, \omega, \theta$ ? From Vallado

### 2.6 Application: $r$ and $v$ from Orbital Elements

We've seen how to find the orbital elements from the position and velocity vectors, but we often need the reverse process to complete certain astrodynamic studies. We'll call the process RANDV to indicate that we're determining the position and velocity vectors. The overall idea is to determine the position and velocity vectors in the perifocal coordinate system, PQW , and then rotate to the geocentric equatorial system. Although the orbit may not be elliptical, and therefore the PQW system would actually be undefined,
we can elegantly work around this limitation. We can also make the method completely generic through several short, simple substitutions.

First, we must use the semiparameter instead of the semimajor axis. As previously mentioned, the semimajor axis is infinite for the parabola, whereas the semiparameter is defined for all orbits. The second requirement concerns how we treat the auxiliary classical orbital elements for the special cases of circular and equatorial orbits.

Let's begin by finding the position and velocity vectors in the perifocal coordinate system. We've developed and presented these equations previously but show them here coupled with the trajectory equation. Notice the use of the semiparameter to replace dependence on the semimajor axis.

$$
\stackrel{\rightharpoonup}{r}_{P Q W}=\left[\begin{array}{c}
\frac{p \cos (\nu)}{1+e \cos (\nu)}  \tag{2-100}\\
\frac{p \sin (\nu)}{1+e \cos (\nu)} \\
0
\end{array}\right]
$$

An immediate difficulty arises when attempting to define the true anomaly for circular orbits. It turns out that the orbital elements may be temporarily replaced with the alternate elements to provide the necessary values for the calculations. Although you can design a change like this so it's transparent to users, make sure any changes or alternate codings use temporary variables and don't alter the original elements. It's possible to substitute values:

IF Circular Equatorial
let $\omega=0.0, \Omega=0.0$, and $\nu=\lambda_{\text {true }}$
IF Circular Inclined

$$
\begin{equation*}
\text { let } \omega=0.0 \text { and } \nu=u \tag{2-101}
\end{equation*}
$$

The rationale for assigning $\omega$ and $\Omega$ to zero will be clear shortly; however, we haven't violated any assumptions because $\omega$ and $\Omega$ are undefined for circular orbits. Be careful not to return any changed variables in computer applications.

Find the velocity vector by differentiating the position vector:

$$
\stackrel{\rightharpoonup}{v}_{P Q W}=\left[\begin{array}{c}
\dot{r} \cos (\nu)-r \nu \operatorname{SIN}(\nu) \\
\dot{r} \operatorname{SIN}(\nu)+r \nu \operatorname{Cos}(\nu) \\
0
\end{array}\right]
$$

Remembering the geometry from Fig. 1-13, solve Eq. (1-18) as

$$
r \nu=\frac{h}{r}
$$

Now, substitute the definitions of position and angular momentum:

$$
r \dot{\nu}=\frac{\sqrt{\mu p}(1+e \cos (\nu))}{p}=\sqrt{\frac{\mu}{p}}(1+e \cos (\nu))
$$

Using Eq. (1-25) and the equation above, write

$$
\dot{r}=\sqrt{\frac{\mu}{p}}(e \sin (\nu))
$$

Substituting these results into the differentiated vector gives us the final solution:

$$
\vec{v}_{P Q W}=\left[\begin{array}{c}
-\sqrt{\frac{\mu}{p}} \operatorname{Sin}(\nu)  \tag{2-102}\\
\sqrt{\frac{\mu}{p}}(e+\cos (\nu)) \\
0
\end{array}\right]
$$

The next step is to rotate the position and velocity vectors to the geocentric equatorial frame. Although this is relatively easy for standard, elliptical, inclined orbits, we'll need to take certain precautions in order to account for special cases, as described with the true anomaly above. We've discussed two of these special cases; the third is the elliptical equatorial case:

$$
\begin{align*}
& \text { IF Elliptical Equatorial }  \tag{2-103}\\
& \text { set } \Omega=0.0 \text { and } \omega=\tilde{\omega}_{\text {true }}
\end{align*}
$$

The assumptions remain intact because $\Omega$ is undefined for elliptical equatorial orbits.
We can now do the coordinate transformations using Eq. (3-28). We may want to multiply out these operations to reduce trigonometric operations. The rationale for setting certain variables to zero should now be apparent. For the special cases, a zero rotation causes the vector to remain unchanged, whereas a desired angular value causes a change.

## Implementing RANDV

Computational efficiency results from assigning the trigonometric terms $[\operatorname{Sin}(\nu)$, $\cos (\nu)]$ and $(\mu / p)$ to temporary variables. This saves many transcendental operations and requires very little extra work. There are also some savings in treating special-case orbits if we reuse the same rotation matrices, but there may be some redundancy in special cases.

As with the ELORB algorithm, we may run many test cases to verify the routine. Because RANDV is simply designed to be a mirror calculation of the ELORB routine, we can use the same set of test reference data. But we must test several limiting cases. Algorithm 10 summarizes the process.

## ALGORITHM 10: RANDV $\left(p, e, i, \Omega, \omega, \nu\left(u, \lambda_{\text {true }}, \tilde{\omega}_{\text {true }}\right) \Rightarrow \stackrel{\rightharpoonup}{r}_{l J K} \stackrel{\rightharpoonup}{v}_{l J K}\right)$

IF Circular Equatorial
$\operatorname{SET}(\omega, \Omega)=0.0$ and $\nu=\lambda_{\text {true }}$
IF Circular Inclined
$\operatorname{SET} \omega=0.0$ and $\nu=u$
IF Elliptical Equatorial

$$
\operatorname{SET} \Omega=0.0 \text { and } \omega=\tilde{\omega}_{\text {true }}
$$

$$
\begin{gathered}
\stackrel{\rightharpoonup}{r}_{P Q W}=\left[\begin{array}{c}
\frac{p \operatorname{COS}(\nu)}{1+e \operatorname{Cos}(\nu)} \\
\frac{p \operatorname{Sin}(\nu)}{1+e \operatorname{COS}(\nu)} \\
0
\end{array}\right] \quad \stackrel{\rightharpoonup}{v}_{P Q W}=\left[\begin{array}{c}
-\sqrt{\frac{\mu}{p}} \operatorname{SIN}(\nu) \\
\sqrt{\frac{\mu}{p}}(e+\operatorname{COS}(\nu)) \\
0
\end{array}\right] \\
\stackrel{\rightharpoonup}{r}_{I / K}=[\operatorname{ROT} 3(-\Omega)][\operatorname{ROT} 1(-i)][\operatorname{ROT} 3(-\omega)] \stackrel{\rightharpoonup}{r}_{P Q W}=\left[\frac{l / K}{P Q W}\right] \stackrel{\rightharpoonup}{r}_{P Q W} \\
\stackrel{\rightharpoonup}{v}_{l / K}=[\operatorname{ROT} 3(-\Omega)][\operatorname{ROT} 1(-i)][\operatorname{ROT} 3(-\omega)] \vec{v}_{P Q W}=\left[\frac{I J K}{P Q W} \vec{v}_{P Q W}\right. \\
{\left[\frac{I J K}{P Q W}\right]=\left[\begin{array}{cc}
\cos (\Omega) \cos (\omega)-\sin (\Omega) \sin (\omega) \cos (i) \\
\sin (\Omega) \cos (\omega)+\cos (\Omega) \operatorname{cin}(\omega) \cos (i) & -\sin (\Omega) \sin (\omega)+\cos (\Omega) \cos (\omega) \cos (i) \\
\sin (\omega) \sin (i) & -\cos (\Omega) \sin (i) \\
\cos (\omega) \sin (i) & \cos (i)
\end{array}\right]}
\end{gathered}
$$

## An example demonstrates the technique.

## Vxample 2-6. Finding Position and Velocity Vectors (RANDV Test Case).

GIVEN: $\quad p=11,067.790 \mathrm{~km}=1.73527 \mathrm{ER}, e=0.83285, i=87.87^{\circ}, \Omega=227.89^{\circ}$, $\omega=53.38^{\circ}, \nu=92.335^{\circ}$
FIND: $\quad \stackrel{\rightharpoonup}{r}_{I J K} \stackrel{\rightharpoonup}{v}_{I J K}$
We have to change the rotation angles if we're using special orbits (equatorial or circular), but this orbit doesn't have special cases. From the given information, form the PQW position and velocity vectors:

$$
\stackrel{r}{P Q W}^{r_{P Q}}\left[\begin{array}{c}
\frac{p \cos (v)}{1+e \cos (\nu)} \\
\frac{p \sin (\nu)}{1+e \cos (\nu)} \\
0
\end{array}\right]=\left[\begin{array}{c}
\frac{1.73527 \cos (92.336)^{\circ}}{1+0.83284 \cos (92.336)^{\circ}} \\
\frac{1.73527 \sin (92.336)}{1+0.83284 \cos (92.336)^{\circ}} \\
0
\end{array}\right]=\left[\begin{array}{c}
-0.0731867 \\
1.7947339 \\
0
\end{array}\right] \mathrm{ER}
$$

$$
\stackrel{\rightharpoonup}{v}_{P Q W}=\left[\begin{array}{c}
-\sqrt{\frac{\mu}{p}} \sin (v) \\
\sqrt{\frac{\mu}{p}}(e+\cos (v)) \\
0
\end{array}\right]=\left[\begin{array}{c}
-\sqrt{\frac{1}{1.73527}} \sin (92.336) \\
\sqrt{\frac{1}{1.73527}}(0.83284+\cos (92.336)) \\
0
\end{array}\right]=\left[\begin{array}{c}
-0.7584998 \\
0.6013136 \\
0
\end{array}\right] \frac{\mathrm{ER}}{\mathrm{TU}}
$$

Rotate these vectors to the geocentric equatorial system using the following rotation matrices:

$$
\begin{aligned}
& \stackrel{\rightharpoonup}{r}_{I J K}=[\operatorname{ROT} 3(-\Omega)][\operatorname{ROT} 1(-i)][\operatorname{ROT} 3(-\omega)] \vec{r}_{P Q W} \\
& \stackrel{\rightharpoonup}{v}_{J J K}=[\operatorname{ROT} 3(-\Omega)][\operatorname{ROT} 1(-i)][\operatorname{ROT} 3(-\omega)] \vec{v}_{P Q W}
\end{aligned}
$$

Or, use the expanded matrix with a computer to do the many trigonometric operations, which result in the transformation matrix

$$
\left[\frac{l J K}{P Q W}\right]=\left[\begin{array}{rrrrr}
-0.37773647 & 0.55459739 & -0.74144244 \\
-0.46253821 & 0.58067014 & 0.669 & 98552 \\
0.80210571 & 0.59602342 & 0.037 & 182 & 20
\end{array}\right]
$$

Finally, multiply each vector to apply the transformation:

$$
\begin{aligned}
& \stackrel{\rightharpoonup}{r}_{I J K}=\left[\frac{U S K}{P Q W}\right]^{\stackrel{\rightharpoonup}{r}}{ }_{P Q W}=\left[\begin{array}{ccc}
-0.37773647 & 0.55459739 & -0.74144244 \\
-0.46253821 & 0.58067014 & 0.66998552 \\
0.80210571 & 0.59602342 & 0.03718220
\end{array}\right]\left[\begin{array}{c}
-0.0731867 \\
1.7947339 \\
0
\end{array}\right] \\
& =\left[\begin{array}{l}
1.023 \\
1.076 \\
1.011
\end{array}\right] \mathrm{ER}=\left[\begin{array}{l}
6524.834 \\
6862.875 \\
6448.296
\end{array}\right] \mathrm{km} \\
& \stackrel{\rightharpoonup}{v}_{I J K}=\left[\frac{U K K}{P Q W}\right] \stackrel{\rightharpoonup}{v}_{P Q W}=\left[\begin{array}{ccc}
-0.37773647 & 0.55459739 & -0.74144244 \\
-0.46253821 & 0.58067014 & 0.66998552 \\
0.80210571 & 0.59602342 & 0.03718220
\end{array}\right]\left[\begin{array}{c}
-0.7584998 \\
0.6013136 \\
0
\end{array}\right] \\
& =\left[\begin{array}{c}
0.62 \\
0.70
\end{array}\right] \mathrm{ER} / \mathrm{TU}=\left[\begin{array}{r}
4.901320 \\
5.533756 \\
1.976 .341
\end{array}\right] \mathrm{km} / \mathrm{s}
\end{aligned}
$$

## Two-Line Elements (TLE)

```
ISS (ZARYA)
1 25544U 98067A 08264.51782528 -.00002182 00000-0 -11606-4 0 2927
2 25544 51.6416 247.4627 0006703 130.5360 325.0288 15.72125391563537
```


## For monitoring by Norad *

The meaning of this data is as follows:

```
LINE 1:
FIELD COLS
CONTENT
    EXAMPLE
    COLS
    03-07
    08-08
    10-11
    12-14
    15-17
    19-20
    21-32
    34-43
    45-52
    54-61
    63-63
    65-68
    69-69
```


## Line numbe

```
1
Satellite number 25544
1 01-01
03-07
Classification (U=Unclassified)
U
International Designator (Last two digits of launch year)
98
10-11
International Designator (Launch number of the year)
067
International Designator (Piece of the launch)
Epoch Year (Last two digits of year)
A
08
19-20
21-32
34-43
45-52
54-61
12 63-63
First Time Derivative of the Mean Motion
264.51782528
\(-.00002182\)
Second Time Derivative of Mean Motion (decimal point assumed)
00000-0
```


## BSTAR drag term (decimal point assumed)

```
-11606-4
The number 0 (Originally this should have been "Ephemeris type")
14 69-69 Checksum (Modulo 10)
292
7
```

```
LINE 2:
FIELD COLS
    01-01
    03-07
    09-16
    18-25
    27-33
    35-42
    44-51
    53-63
    64-68
    69-69
```

CONTENT
Line number
Satellite number
Inclination [Degrees]
Right Ascension of the Ascending Node [Degrees]
Eccentricity (decimal point assumed)
Argument of Perigee [Degrees]
Mean Anomaly [Degrees]
Mean Motion [Revs per day]
Revolution number at epoch [Revs]
Checksum (Modulo 10)

## EXAMPLE

## 2

25544
51.6416
247.4627

0006703
130.5360
325.0288
15.72125391

56353
7

## Celestrak: Update TLE



## Celestrak: ISS, February 24, 2009



## Lost ISS Toolbag



## Heidemarie Stefanyshyn-Piper

From Wikipedia, the free encyclopedia
Heidemarie Martha Stefanyshyn-Piper (born on February 7, 1963 experienced salvage officer. Her major salvage projects include de-s Peruvian submarine Pacocha.

Stefanyshyn-Piper has received numerous honors and awards, suct 115 and STS-126, during which she completed five spacewalks tota
$\quad$ Contents [hide]
1 Early life and education
2 Military career
3 NASA career
3.1 STS-115 - Atlantis (September 9-21, 2006)
3.2 NEEMO 12 (May 7-18, 2007)
3.3 STS-126 - Endeavour (November 14-30, 2008)
$\quad$ 3.3.1 Lost tool bag during spacewalk
4 Retirement from NASA
5 Commanding the NSWCCD
6 References
7 External links

## Celestrak: IRIDIUM 33, February 24, 2009



# Astrodynamics <br> (AEROOO24) 

## 3B. The Orbit in Space

## Gaëtan Kerschen

## Space Structures \& Systems Lab (S3L)

